

The large-time asymptotics for the modified Camassa-Holm equation on a non-zero background

Iryna Karpenko^{†,‡} and Dmitry Shepelsky[†]

[†] B. Verkin Institute for Low Temperature Physics and Engineering of NAS of Ukraine

[‡] University of Vienna

Cauchy problem for the mCH equation

The modified Camassa-Holm (mCH) equation:

$$m_t + \left((u^2 - u_x^2)m \right)_x = 0, \quad m = u - u_{xx}.$$

Cauchy problem:

$$u(x, 0) = u_0(x), \quad -\infty < x < \infty,$$

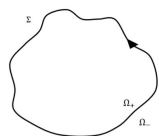
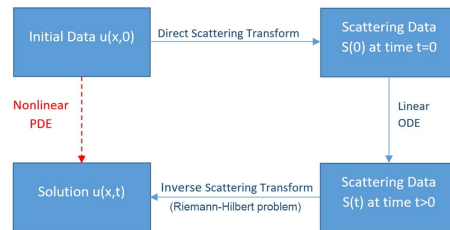
where $u_0(x) \rightarrow 1$ as $|x| \rightarrow \infty$.

- Introducing \tilde{u} such that $u(x, t) = \tilde{u}(x - t, t) + 1$, the mCH equation becomes

$$\tilde{m}_t + \left((\tilde{u}^2 - \tilde{u}_x^2 + 2\tilde{u})\tilde{m} \right)_x = 0, \quad \tilde{m} = \tilde{u} - \tilde{u}_{xx} + 1$$

on the zero background ($\tilde{u} \rightarrow 0$ as $|x| \rightarrow \infty$).

Approach: Riemann-Hilbert problem



Riemann-Hilbert problem: boundary value problem in complex analysis

Given contour $\Sigma \in \mathbb{C}$ and "jump function" $G(s) : \Sigma \rightarrow \mathbb{C}^{n \times n}$, find a function $M(z) : \mathbb{C} \setminus \Sigma \rightarrow \mathbb{C}^{n \times n}$ such that:

- $M(z)$ is analytic in $\mathbb{C} \setminus \Sigma$;
- $M_+(s) = M_-(s)G(s)$, $s \in \Sigma$;
- $M(\infty) = I$.

- In applications, the jump matrix G depends on **parameters**; then the solution also depends on these **parameters**; e.g., $G(s; x, t) \rightarrow M(z; x, t)$

Inverse scattering transform method in the RHP form

- using
 - the **Lax pair** associated to the mCH equation
- construct
 - a multiplicative matrix **Riemann-Hilbert problem (RHP)**
- obtain
 - a representation of the solution $u(x, t)$ of the mCH equation in terms of the solution $M(x, t; \cdot)$ of the associated RHP
- obtain
 - the long-time asymptotics of $u(x, t)$ via the Deift-Zhou **nonlinear steepest descent** method.

mCH as integrable equation

- This equation is the compatibility condition for the Lax pair equations:

$$\begin{cases} \Phi_x(x, t, \lambda) = U(x, t, \lambda)\Phi(x, t, \lambda) \\ \Phi_t(x, t, \lambda) = V(x, t, \lambda)\Phi(x, t, \lambda) \end{cases},$$

where U and V are 2×2 matrices depending on the

spectral parameter λ : $U = \frac{1}{2} \begin{pmatrix} -1 & \lambda \tilde{m} \\ -\lambda \tilde{m} & 1 \end{pmatrix}$;

$$V = \begin{pmatrix} \lambda^{-2} + \frac{(\tilde{u}^2 - \tilde{u}_x^2 + 2\tilde{u})}{2} & -\lambda^{-1}(\tilde{u} - \tilde{u}_x + 1) - \frac{\lambda(\tilde{u}^2 - \tilde{u}_x^2 + 2\tilde{u})\tilde{m}}{2} \\ \lambda^{-1}(\tilde{u} + \tilde{u}_x + 1) + \frac{\lambda(\tilde{u}^2 - \tilde{u}_x^2 + 2\tilde{u})\tilde{m}}{2} & -\lambda^{-2} - \frac{(\tilde{u}^2 - \tilde{u}_x^2 + 2\tilde{u})}{2} \end{pmatrix}.$$

- In order to control large λ behaviour, we introduce a new spatial variable

$$y(x, t) = x - \int_x^{+\infty} (\tilde{m}(\xi, t) - 1) d\xi.$$

Obtaining $\tilde{u}(x, t)$

Consider the **solitonless** case assuming that there are no residue conditions.

- Given $u_0(x)$, construct the "reflection coefficient" $r(\mu)$, $\mu \in \mathbb{R}$, by solving the Lax pair equations, whose coefficients are determined in terms of $u_0(x)$.
- Construct the jump matrix $J(y, t, \mu)$, $\mu \in \mathbb{R}$ by

$$J(y, t, \mu) = e^{-p(y, t, \mu)\sigma_3} J_0(\mu) e^{p(y, t, \mu)\sigma_3}$$

(J explicitly depends on y and t) where

$$p(y, t, \mu) = \frac{i(\mu^2 - 1)}{4\mu} \left(-y + \frac{8\mu^2}{(\mu^2 + 1)^2} t \right)$$

and $J_0(\mu)$ is defined by

$$J_0(\mu) = \begin{pmatrix} 1 - r(\mu)r^*(\mu) & r(\mu) \\ -r^*(\mu) & 1 \end{pmatrix}.$$

- Solve the following **RH problem**: Find a piece-wise meromorphic, 2×2 -matrix valued function $M(y, t, \mu)$ satisfying:

- $M_{\pm}(y, t, \mu) = M_{\mp}(y, t, \mu) e^{-p\sigma_3} \begin{pmatrix} 1 - r(\mu)r^*(\mu) & r(\mu) \\ -r^*(\mu) & 1 \end{pmatrix} e^{p\sigma_3}$, $\mu \in \mathbb{R} \setminus \{-1, 1\}$, where $p := p(y, t, \mu) = \frac{i(\mu^2 - 1)}{4\mu} \left(-y + \frac{8\mu^2}{(\mu^2 + 1)^2} t \right)$ and $\sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$
- $M \rightarrow I$ as $\mu \rightarrow \infty$
- $M_{\pm}(y, t, \mu) = \frac{1}{2(\mu-1)} \alpha_{\pm}(y, t) \begin{pmatrix} -\zeta & 1 \\ \zeta & -1 \end{pmatrix} + O(1)$ as $\mu \rightarrow 1$
- $M_{\pm}(y, t, \mu) = \frac{1}{2(\mu+1)} \alpha_{\pm}(y, t) \begin{pmatrix} -\zeta & -1 \\ \zeta & 1 \end{pmatrix} + O(1)$ as $\mu \rightarrow -1$
- $M(\cdot, \cdot, \frac{1}{\mu}) = \tilde{M}(\cdot, \cdot, \bar{\mu}) = \sigma_3 M(\cdot, \cdot, -\mu) \sigma_3 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} M(\cdot, \cdot, \mu) \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$,

- Expand $M(y, t, \mu)$ in the neighborhood of i in order to obtain $a(y, t)$, $a_2(y, t)$, $a_3(y, t)$:

$$M(y, t, \mu) = \begin{pmatrix} a(y, t) & 0 \\ 0 & \frac{1}{a(y, t)} \end{pmatrix} + \begin{pmatrix} 0 & a_2(y, t) \\ a_3(y, t) & 0 \end{pmatrix} (\mu - i) + O((\mu - i)^2)$$

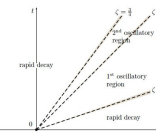
Then \tilde{u} is given in a parametric form:

- $\tilde{u}(y, t) = -\frac{a_3(y, t)}{a(y, t)} - a_2(y, t)a(y, t)$
- $\tilde{u}_x(y, t) = \frac{a_3(y, t)}{a(y, t)} - a_2(y, t)a(y, t)$
- $x(y, t) = y + 2 \ln a(y, t)$

Asymptotics

Question

How does $u(x, t)$ behave for large t ?

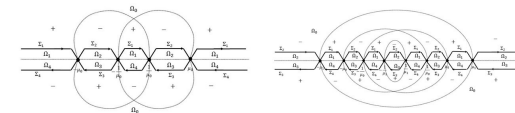


Answer

- four sectors** in the (x, t) half-plane where $u(x, t)$ behaves differently for large t , depending on the magnitude of $\zeta = \frac{x}{t}$

Signature table:

Sign of $\text{Im} \frac{p(y, t, \mu)}{i\mu} = \frac{1}{4} \left(\mu - \frac{1}{\mu} \right) \left(\xi - \frac{2}{(\mu + \frac{1}{\mu})^2} \right)$ for various $\xi = \frac{y}{t}$; μ near \mathbb{R}



$1 < \zeta < 3$

$\frac{3}{4} < \zeta < 1$

$1 < \frac{x}{t} < 3$

In the **solitonless** case, the solution $u(x, t)$ of the Cauchy problem for the mCH equation has the following large-time asymptotics in the sector of the (x, t) half-plane defined by $1 < \zeta := \frac{x}{t} < 3$:

$$u(x, t) = 1 + \frac{C_1(\zeta - 1)}{\sqrt{t}} \cos \left\{ C_2(\zeta - 1)t + C_3(\zeta - 1) \ln t + \tilde{C}_4(\zeta - 1) \right\} + o(t^{-1/2}),$$

where C_i are functions of ζ specified in terms of the scattering data, which in turn are uniquely specified by the initial data.

The error term is uniform in any sector $1 + \varepsilon < \zeta < 3 - \varepsilon$ where ε is a small positive number.

$\frac{3}{4} < \frac{x}{t} < 1$

In the **solitonless** case, the solution $u(x, t)$ of the Cauchy problem for the mCH equation has the following large-time asymptotics along the rays $\frac{x}{t} := \zeta$ in the sector of the (x, t) half-plane defined by $\frac{3}{4} < \zeta < 1$:

$$u(x, t) = 1 + \sum_{j=0,1} \frac{C_1^{(j)}(\zeta - 1)}{\sqrt{t}} \cos \left(C_2^{(j)}(\zeta - 1)t + C_3^{(j)}(\zeta - 1) \ln t + \tilde{C}_4^{(j)}(\zeta - 1) \right) + o(t^{-1/2}),$$

where $C_i^{(j)}$ are functions of ζ specified in terms of the scattering data, which in turn are uniquely specified by the initial data.

The error term is uniform in any sector $\frac{3}{4} + \varepsilon < \zeta < 1 - \varepsilon$ where ε is small and positive.

References

- Anne Boutet de Monvel, Iryna Karpenko, Dmitry Shepelsky, The modified Camassa-Holm equation on a nonzero background: large-time asymptotics for the Cauchy problem, preprint: arXiv:2011.13235
- Anne Boutet de Monvel, Iryna Karpenko, Dmitry Shepelsky, A Riemann-Hilbert approach to the modified Camassa-Holm equation with nonzero boundary conditions, Journal of Mathematical Physics, Volume 61, Issue 3 (2020), 031504.